



Elastic co-tunnelling in SET's - and other s-qubit issues -

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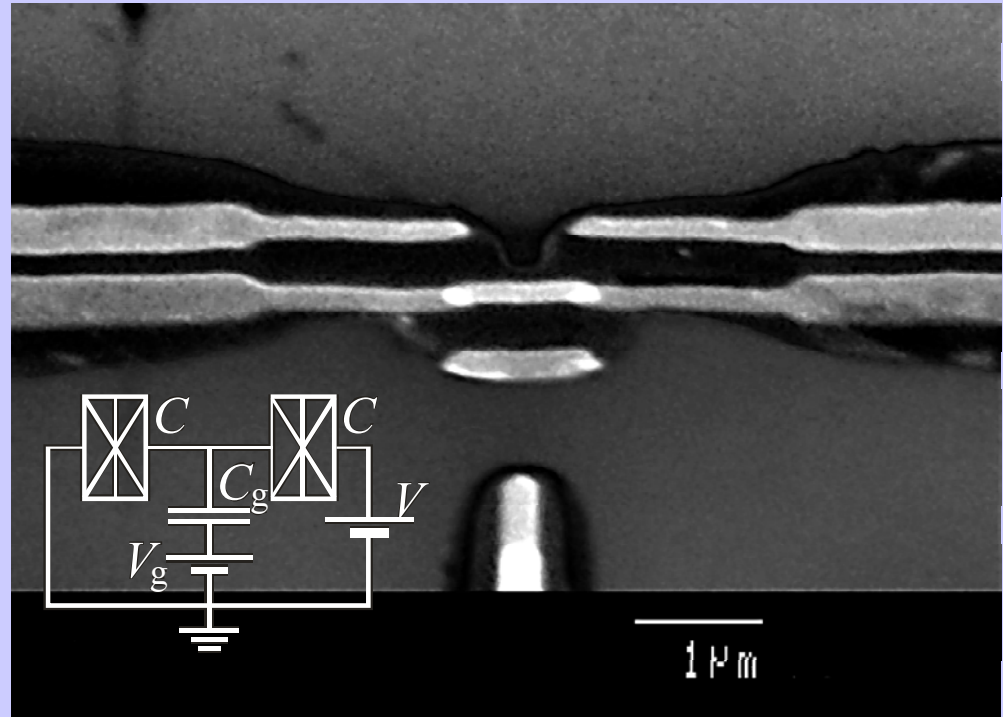
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Nb as a construction material for nano-devices:

Bulk properties/Material	Aluminum	Niobium
Critical temperature	1.175 K (0.1 meV)	9.26 K (0.8 meV)
Superconducting gap	2.07 K (0.18 meV)	16.33 K (1.4 meV)
Critical magnetic field	0.0105 T	0.2060 T

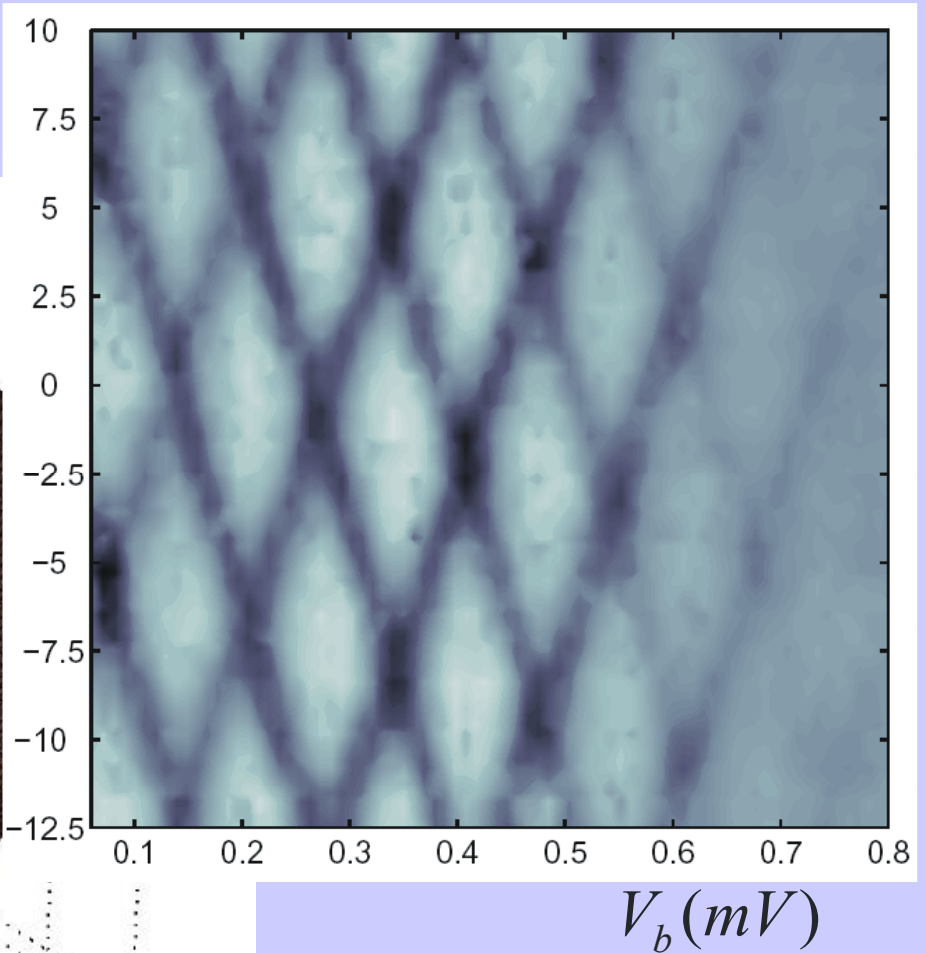
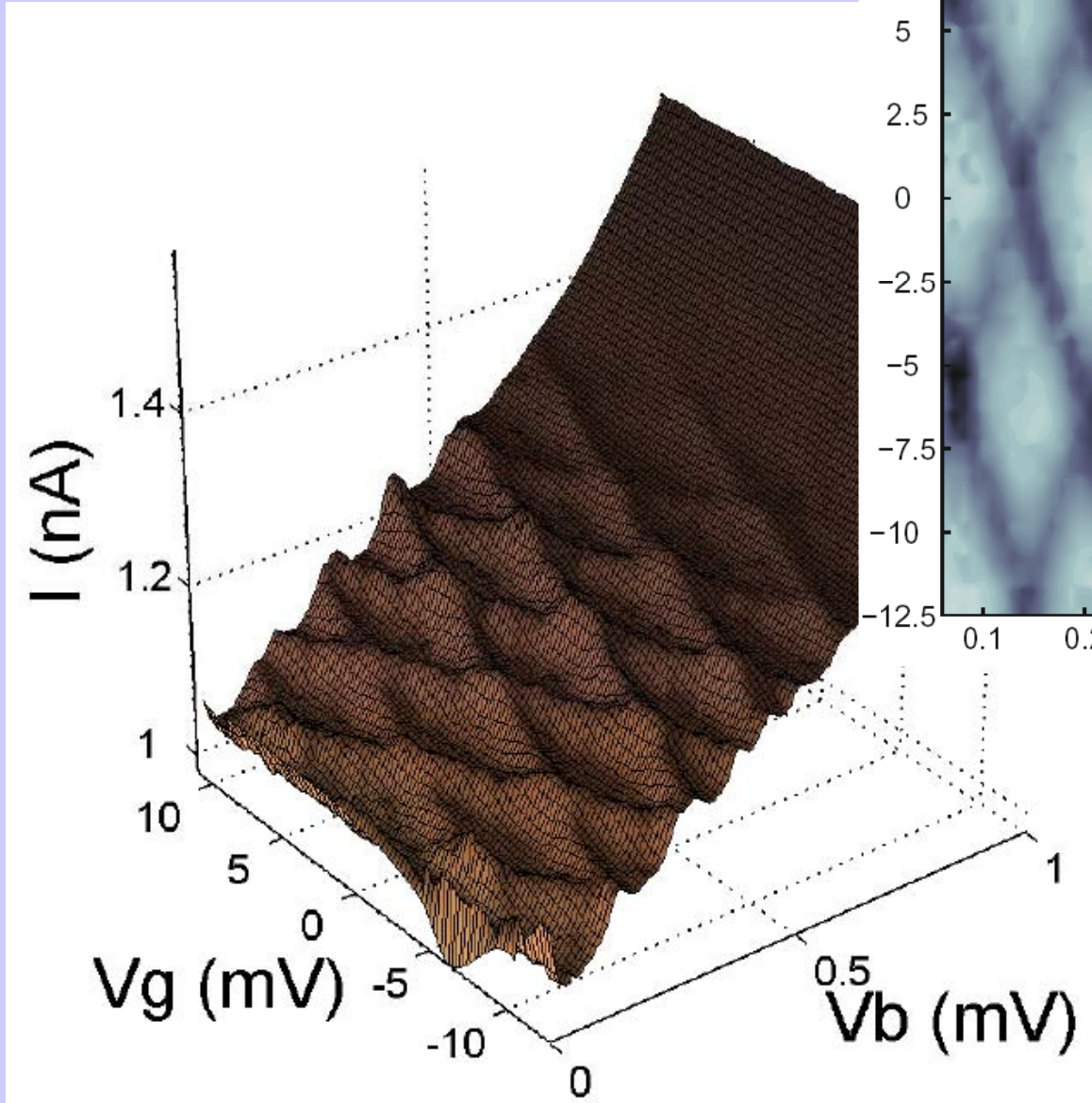
$$\mu_0 = 4\pi \times 10^{-7} \text{ H / m}$$

Single-electron transistor Al/Nb/Al:



Cooper-pair resonances:

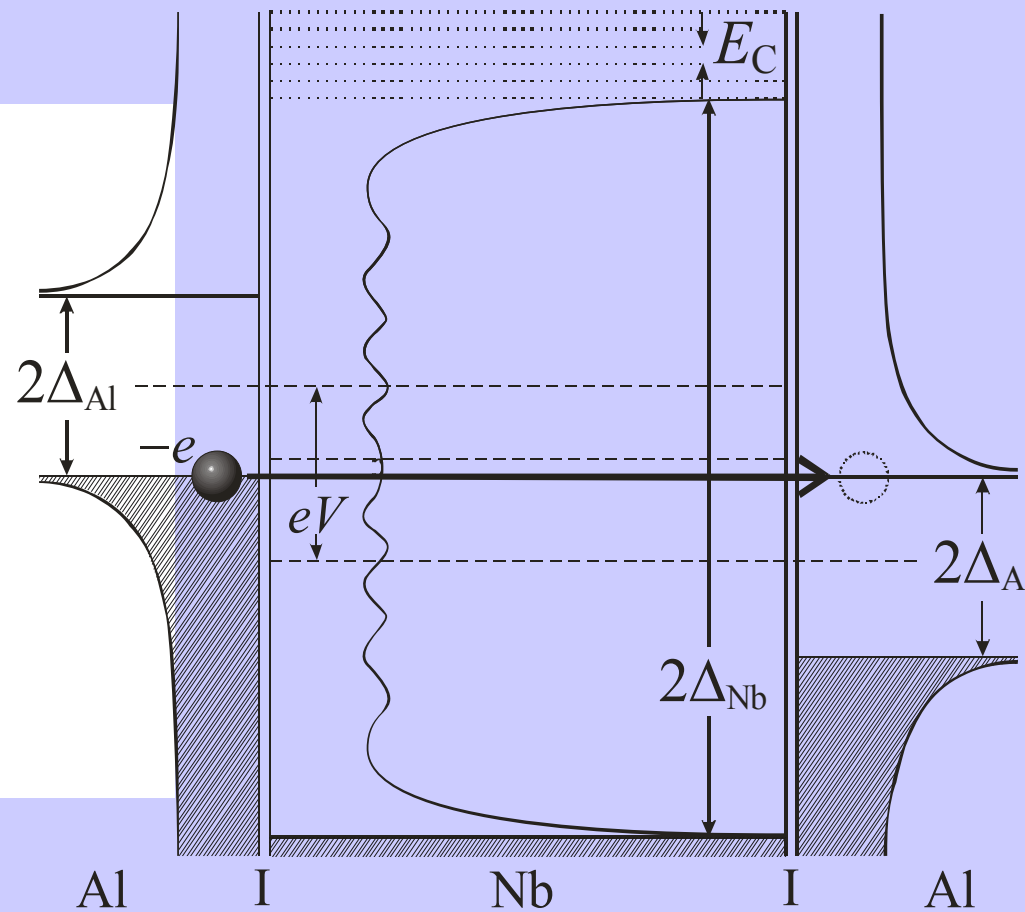
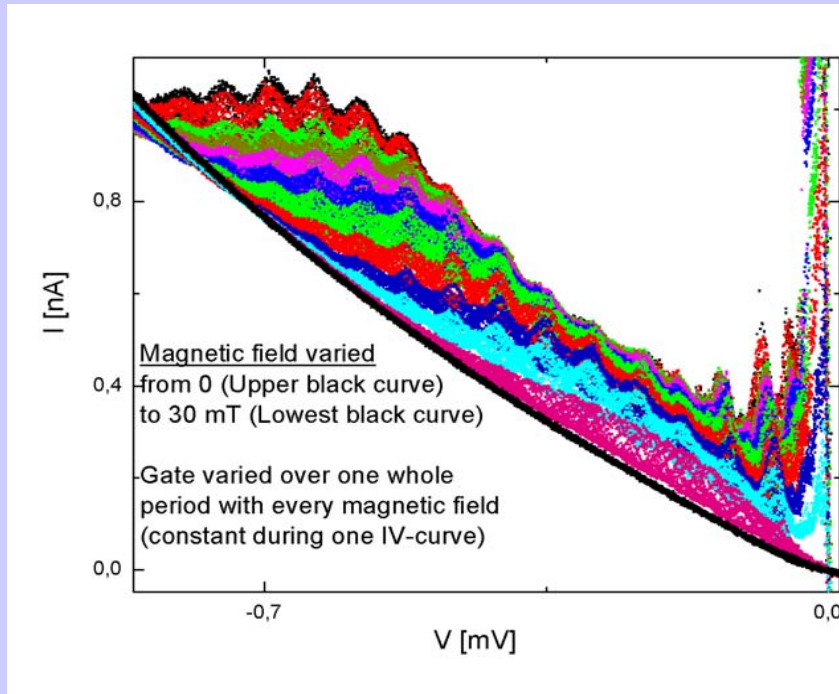
V_g (mV)



V_b (mV)

$$V = \frac{1}{C_\Sigma} (2e + 2C_g V_g + 2Q_0)$$

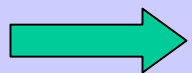
Measurements in magnetic field:



SIMPLEST INTERPRETATION:
 assume that there are "leaky" electronic states inside the gap of Nb.
 Then the background current would be given by sequential tunneling through these states.

$$\Gamma_{LI}(n) = \frac{1}{e^2 R_{LI}} \iint d\varepsilon d\varepsilon' N_L(\varepsilon) N_I(\varepsilon') f_L(\varepsilon) [1 - f_L(\varepsilon')] \delta[E_{CH}(n) + \varepsilon' - \varepsilon]$$

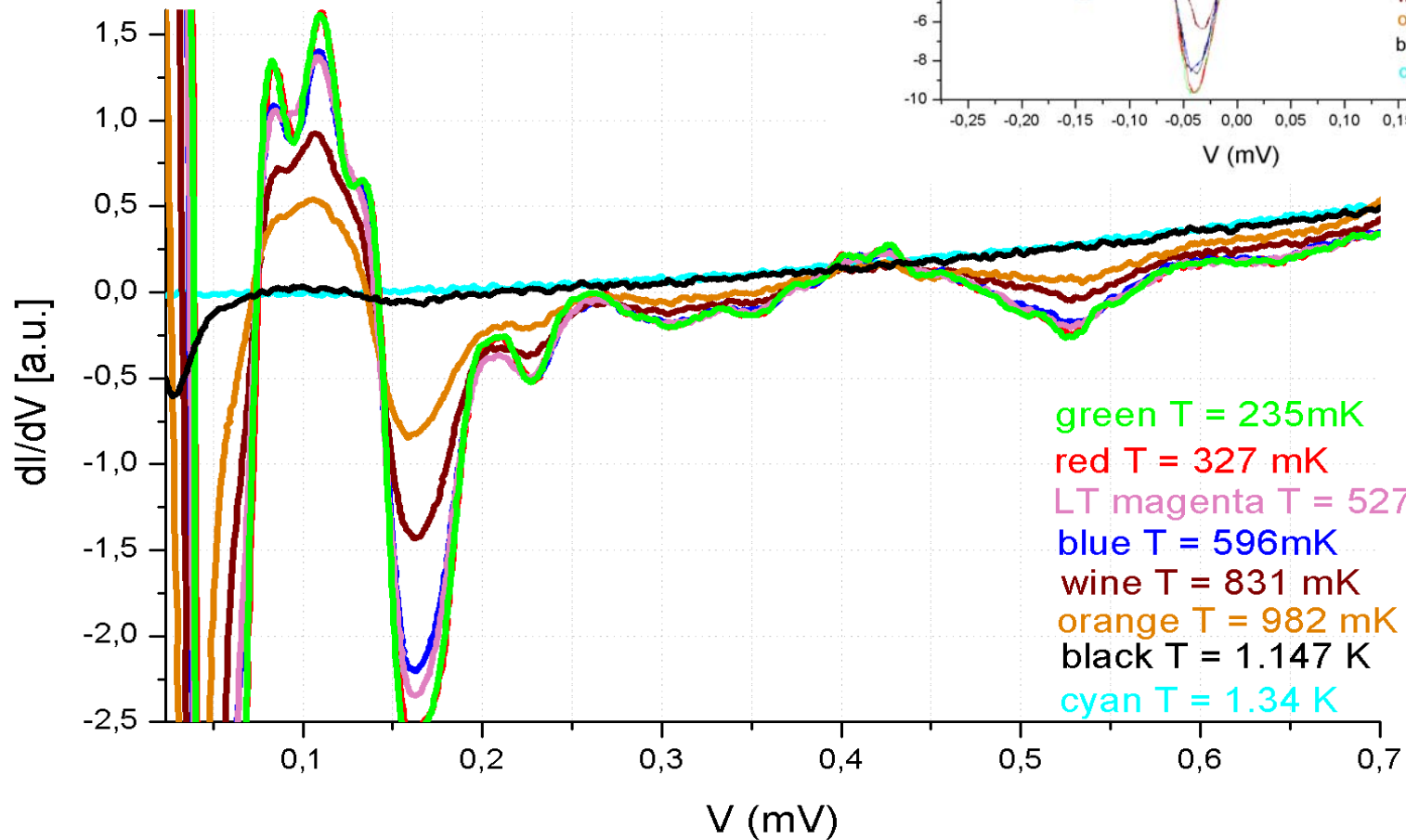
$$N_L(\varepsilon) = \Theta(|\varepsilon| - \Delta_{Al}) \frac{|\varepsilon|}{\sqrt{\varepsilon^2 - \Delta_{Al}^2}}$$



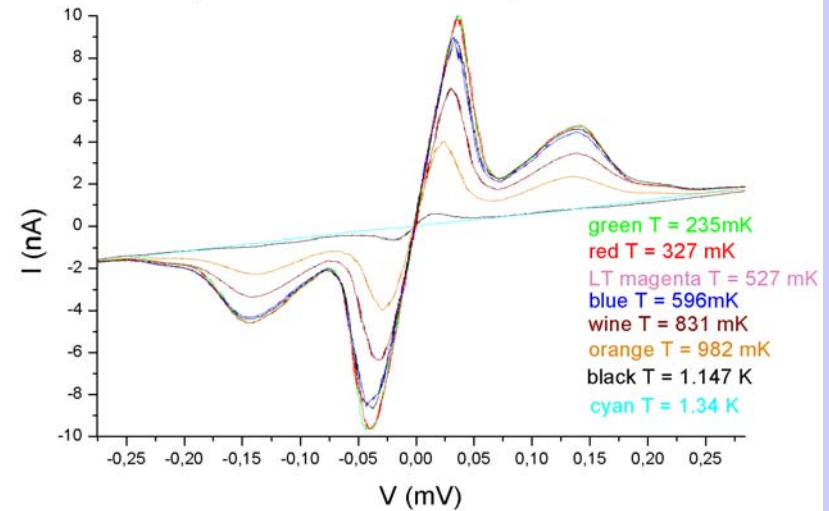
then use the "orthodox" theory, etc.

"Leak" currents in a single AlNb junction:

Conductances at different temperatures



Supercurrents at various temperatures



The theory of elastic co-tunneling:

1. Normal-metal SET's: elastic co-tunneling through the Coulomb barrier.

D. V. Averin and Yu. V. Nazarov, PRL 65 2446 (1990)

2. Superconducting SET's: elastic co-tunneling through the gap of the island and the Coulomb barrier

$$E_1 = \frac{e}{C_2} \left[e \left(n + \frac{1}{2} \right) - V \left(C_2 + \frac{C_g}{2} \right) + C_g V_g \right]$$

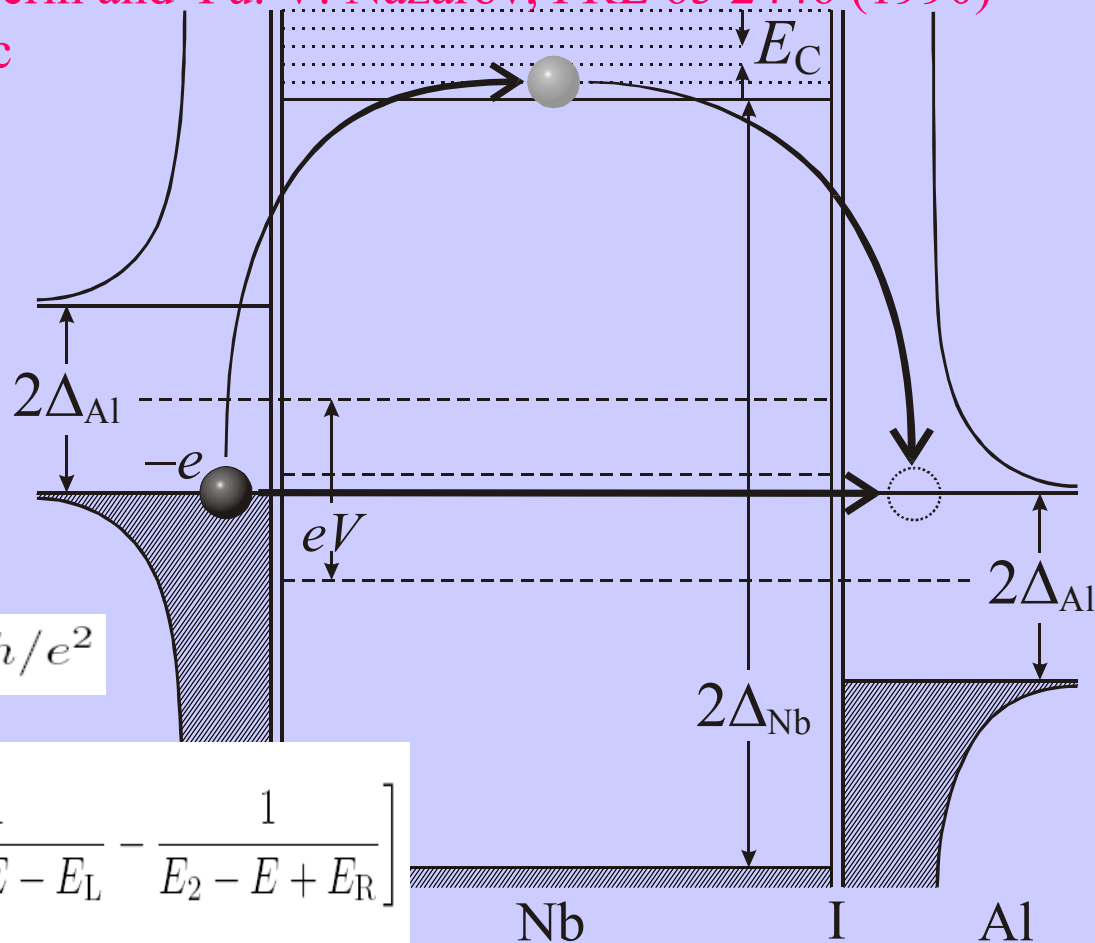
$$E_2 = \frac{e}{C_2} \left[-e \left(n - \frac{1}{2} \right) - V \left(C_1 + \frac{C_g}{2} \right) - C_g V_g \right]$$

$$R_{\text{eff}} = R_{T,1} R_{T,2} / R_K$$

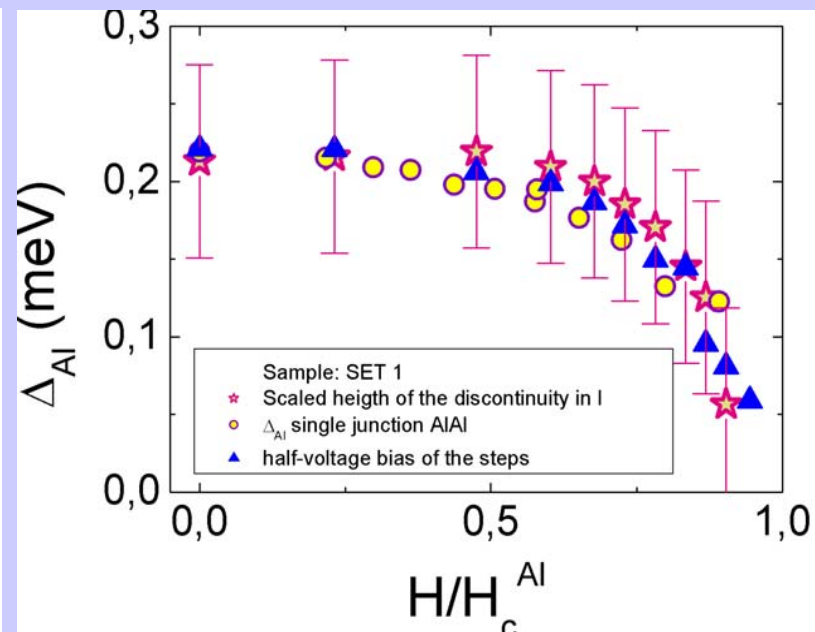
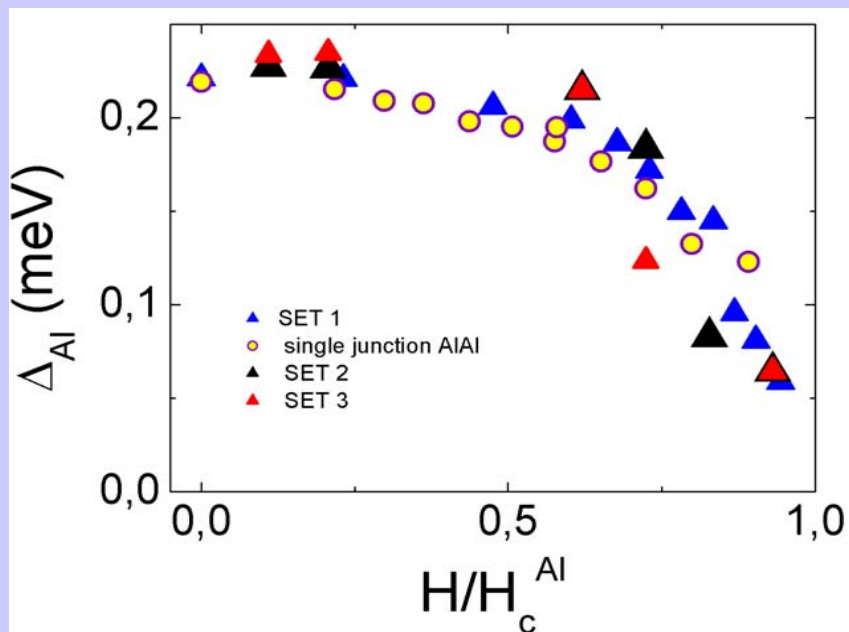
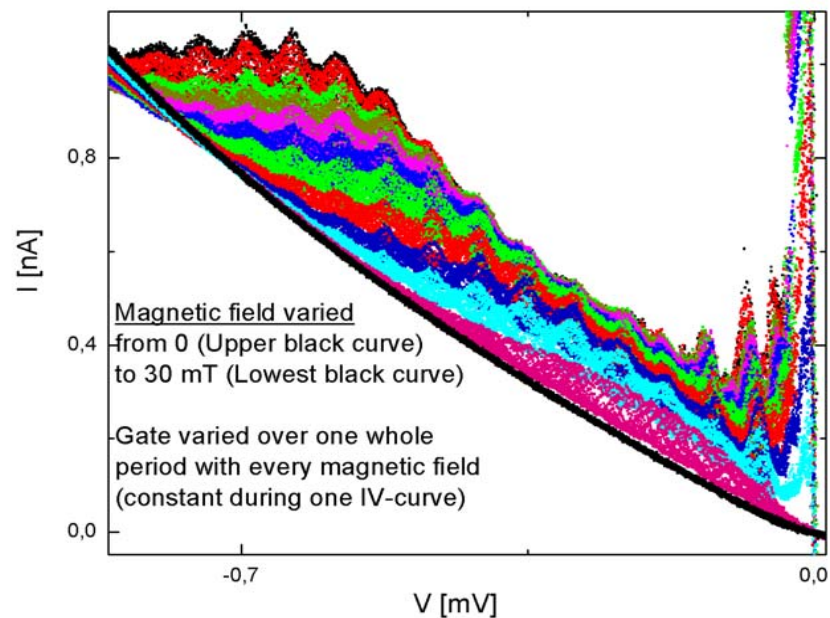
$$R_K = h/e^2$$

$$T_{\text{eff}}(E_L, E_R) = 2\pi \int_{\Delta_{\text{Nb}}}^{\infty} dE \frac{E}{\sqrt{E^2 - \Delta_{\text{Nb}}^2}} \left[\frac{1}{E_1 + E - E_L} - \frac{1}{E_2 - E + E_R} \right]$$

$$I_{\text{el}} = \frac{1}{e R_{\text{eff}}} \int_{\Delta_{\text{Al}}}^{\infty} dE_L \int_{\Delta_{\text{Al}}}^{\infty} dE_R \frac{E_L E_R}{\sqrt{E_L^2 - \Delta_{\text{Al}}^2} \sqrt{E_R^2 - \Delta_{\text{Al}}^2}} T_{\text{eff}}^2(E_L, E_R) \delta(eV - E_L - E_R)$$



Background current: -behaviour in magnetic field -comparison with the elastic co-tunnelling theory

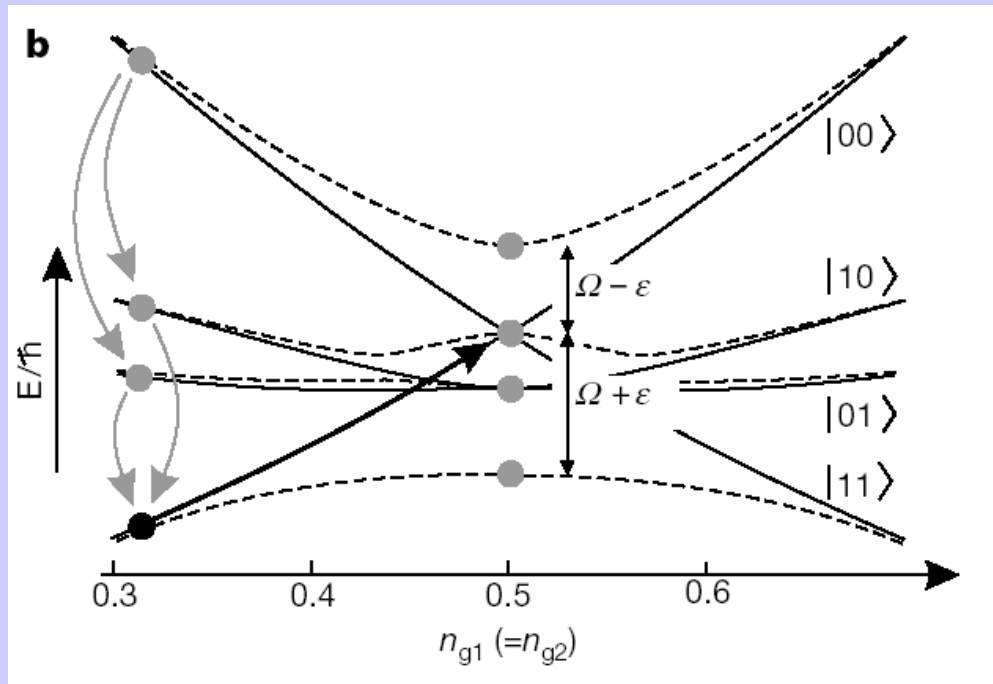
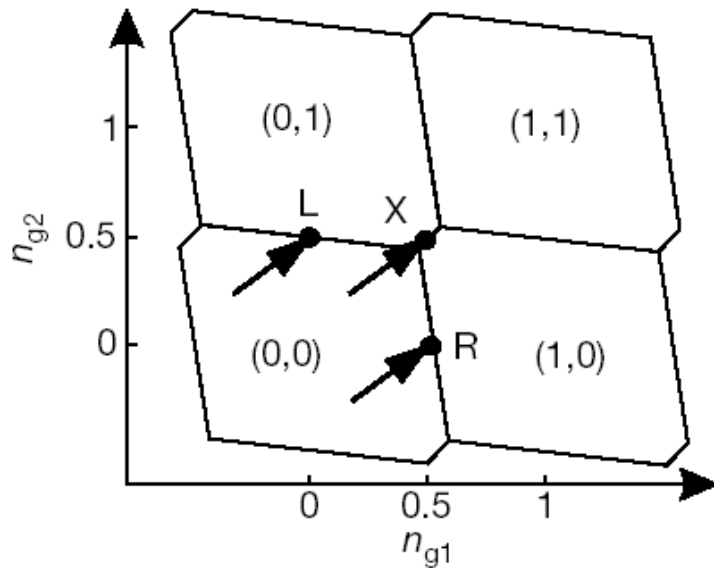


TWO QUBITS ISSUES:-

quantum manipulations with r.f.
pulses -

Generic form for the
2-qubit Hamiltonian
(with interaction):

$$H = \begin{bmatrix} E_{00} & -\frac{1}{2}E_{J1} & -\frac{1}{2}E_{J2} & 0 \\ -\frac{1}{2}E_{J1} & E_{10} & 0 & -\frac{1}{2}E_{J2} \\ -\frac{1}{2}E_{J2} & 0 & E_{01} & -\frac{1}{2}E_{J1} \\ 0 & -\frac{1}{2}E_{J2} & -\frac{1}{2}E_{J1} & E_{11} \end{bmatrix} + \text{measurement part}$$



Yu. A. Pashkin *et al.*, Nature 421, 823(2003).

$$\varepsilon = \frac{1}{2\hbar} \sqrt{(E_{J2} - E_{J1})^2 + \left(\frac{E_m}{2}\right)^2}$$

$$\Omega = \frac{1}{2\hbar} \sqrt{(E_{J1} + E_{J2})^2 + \left(\frac{E_m}{2}\right)^2}$$

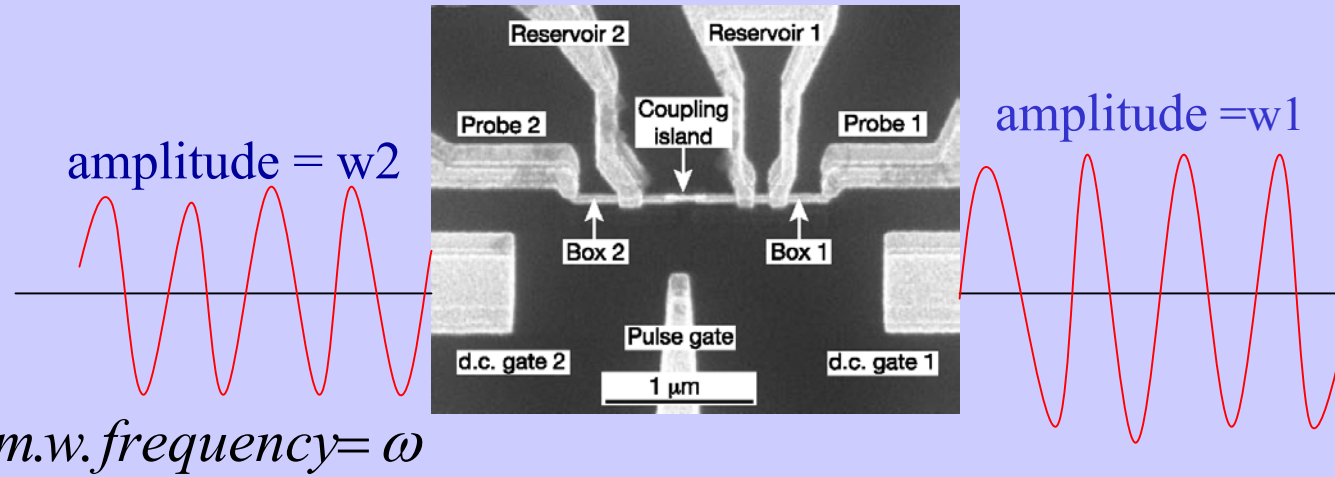
2-qubit Rabi pulses:

Detunings:

$$\delta = \omega - (\Omega - \varepsilon)$$

$$\Delta = \omega - (\Omega + \varepsilon)$$

R.W.A. in a
four-level system!



$$\theta = -\arctan \frac{E_m}{E_{J1} + E_{J2}}$$

$$\chi = -\arctan \frac{E_m}{E_{J2} - E_{J1}}$$

$$H = \begin{bmatrix} \frac{\hbar}{2}(\delta + \Delta) & M_{1+} & M_{2+} & 0 \\ M_{1+} & \frac{\hbar}{2}(-\delta + \Delta) & 0 & M_{2-} \\ M_{2+} & 0 & \frac{\hbar}{2}(\delta - \Delta) & M_{1-} \\ 0 & M_{2-} & M_{1-} & \frac{\hbar}{2}(-\delta - \Delta) \end{bmatrix}$$

2-qubit operations:

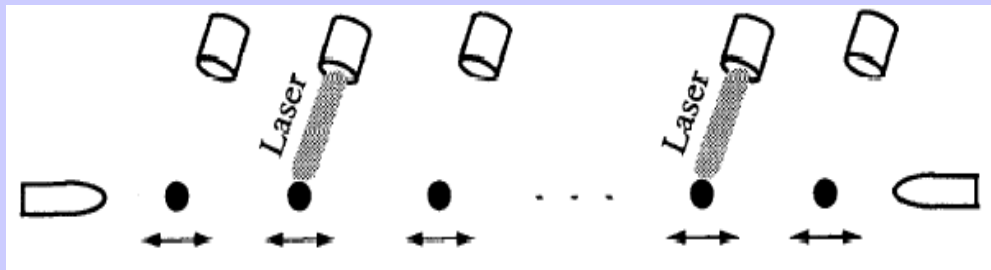
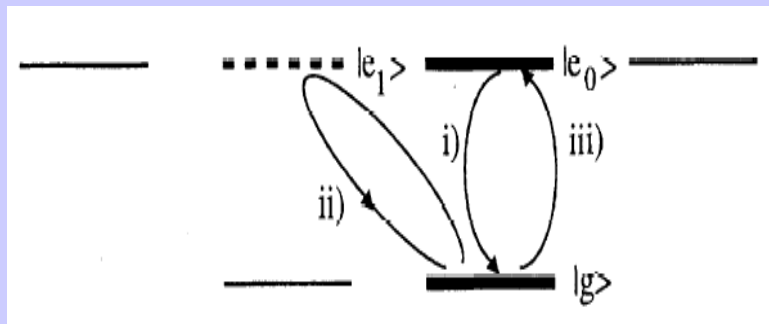
the "almost" theorem of
S. Lloyd, PRL 75, 346 (1995).

$$M_{1\pm} = \frac{E_{c1} w_1}{2} \cos \frac{\theta - \chi}{2} \pm \frac{E_{c2} w_2}{2} \sin \frac{\theta + \chi}{2}$$

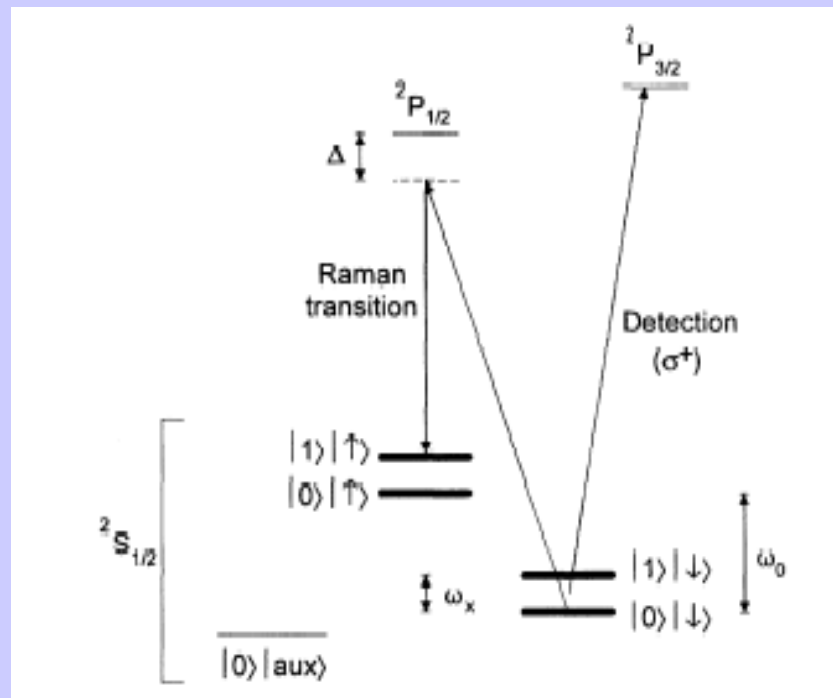
$$M_{2\pm} = \frac{E_{c2} w_2}{2} \cos \frac{\theta + \chi}{2} \pm \frac{E_{c1} w_1}{2} \sin \frac{\theta - \chi}{2}$$

Comparison: trapped ions as qubits

Theory – Cirac&Zoller: PRL 74, 4091 (1995)



Experiment - NIST: PRL 75, 4714 (1995)



$$\hat{H}_{n,q} = \frac{\eta}{\sqrt{N}} \frac{\Omega}{2} [|e_q\rangle_n \langle g| a e^{-i\phi} + |g\rangle_n \langle e_q| a^\dagger e^{i\phi}]$$

$$\eta = [\hbar k_\theta^2 / (2M \nu_x)]^{1/2} = \text{Lamb-Dicke parameter}$$

$$H = \frac{\eta\Omega}{2\sqrt{N}} \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & \exp(-i\phi) & 0 \\ 0 & \exp(i\phi) & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}$$